

ORBITS OF STARS IN A MODEL WITH THREE FIRST INTEGRALS

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RESUMEN. Se aplica una forma integrable del potencial de la Galaxia para describir la variación de los pericentros y apocentros de las órbitas de las estrellas de acuerdo con sus velocidades en el entorno solar. Como ejemplo, se presenta un estudio más detallado de la órbita del Sol.

ABSTRACT. An integrable form of the potential of the Galaxy is applied to describe the variation of the pericentres and the apocentres of the orbits of stars in accordance with their velocities in the solar neighbourhood. As an example, a more detailed study of the orbit of the Sun is presented.

Key words: DYNAMICS -- GALAXY-STRUCTURE

. INTRODUCTION

Chandrasekhar (1942) developed a theory for the study of the dynamics of the stellar systems. Among its main consequences, stellar systems are described as integrable systems: the different forms of the galactic potential U per unit mass have sets of first integrals that are in involution. A short description of the formulation of the problem and the derivation of stationary axisymmetric solutions is given in Cortés and Sala (1990) and a more detailed description of time depending axisymmetric solutions in Sala (1990).

Dejonghe and de Zeeuw (1989a) have approximated within a 3% error the potential derived by Aguilar from the model of the Galaxy by Bahcall and Ostriker (1980) by an Eddington type potential

where λ, ϕ, ν are prolate spheroidal coordinates, being the foci of the spheroids $\lambda = \pm a$ and the hyperboloids of revolution placed on the z -axis at $z = (\gamma - \alpha)^{1/2} = \pm 0.88$ kpc with $\gamma = -10^6$ and where F can be found in Dejonghe and de Zeeuw (1989a).

I. THE ORBITS OF STARS

Having the galactic potential U the Eddington form (1.1) the orbits of stars have the three first integrals

$$\begin{aligned} I_1 &\equiv \Pi^2 + \Phi^2 + Z^2 + 2 \frac{F(\lambda) - F(\nu)}{\lambda - \nu} = \text{constant} \\ I_2 &\equiv \omega \Phi = \text{constant} \\ I_3 &\equiv (z\Pi - \omega Z)^2 + z^2 \Phi^2 + (\gamma - \alpha) Z^2 + \frac{\nu F(\lambda) - \lambda F(\nu)}{\lambda - \nu} = \text{constant} \end{aligned} \quad (2.1)$$

where ω, ϕ, z are the cylindrical coordinates and Π, Φ, Z the physical cylindrical velocity components of a star.

The first integrals (2.1) are in involution, so that, if they are replaced by

$$\begin{aligned}
 H &= \frac{I_1}{2} = \frac{\rho_\lambda^2}{2P^2} + \frac{\rho_\phi^2}{2\omega^2} + \frac{\rho_\nu^2}{2R^2} + \frac{F(\lambda) - F(\nu)}{\lambda - \nu} = \text{constant} \\
 I &= I_2 = \rho_\phi = \text{constant} \\
 J &= \frac{I_3 + I_2^2 - \gamma I_1}{2} = \frac{\nu\rho_\lambda^2}{2P^2} + \frac{(\lambda+\nu+\alpha)\rho_\phi^2}{2\pi^2} + \frac{\lambda\rho_\nu^2}{2R^2} + \frac{\nu F(\lambda) - \lambda F(\nu)}{\lambda - \nu} = \text{constant}
 \end{aligned}$$

where $\rho_\lambda = P^2\lambda'$, $\rho_\phi = \omega^2\phi'$, $\rho_\nu = R^2\nu'$ are the momenta conjugate to the spheroidal coordinates λ, ϕ, z , being its time derivatives denoted by λ', ϕ', ν' and being

$$P^2 = \frac{\lambda - \nu}{4(\lambda+\alpha)(\lambda+\gamma)} \quad R^2 = \frac{\nu - \lambda}{4(\nu+\alpha)(\nu+\gamma)}$$

the equations of the motion of a star in a meridional plane can be expressed

$$\begin{aligned}
 \lambda' &= \left[\frac{8(\lambda+\alpha)(\lambda+\gamma)}{(\lambda - \nu)^2} (\lambda H + \frac{\alpha - \gamma}{2(\lambda+\alpha)} I^2 - J - F(\lambda)) \right]^{1/2} \\
 \nu' &= \left[\frac{8(\nu+\alpha)(\nu+\gamma)}{(\nu - \lambda)^2} (\nu H + \frac{\alpha - \gamma}{2(\nu+\alpha)} I^2 - J - F(\nu)) \right]^{1/2}
 \end{aligned} \quad (2.2)$$

so that they are separable and can be integrated as quadratures (Lynden-Bell 1962). Finally, the motion around the rotation axis will be given by integrating $\phi' = 1/\omega^2$ where $\omega^2 = (\lambda+\alpha)(\nu+\alpha)/(\alpha-\gamma)$ has been obtained from the integration of (2.2). The obtained orbits with non-zero angular momentum are stable short axis tubes limited by spheroids and hyperboloids of revolution (Dejonghe and de Zeeuw 1989b) being both direct and retrograde orbits possible. These limiting coordinate surfaces are derived from the equations $\lambda' = 0$ and $\nu' = 0$ obtained from (2.2) so that a meridional section of a short axis tube is determined by the conditions $\lambda_{\min} \leq \lambda \leq \lambda_{\max}$ and $-\gamma \leq \nu \leq \nu_{\max}$. A more graphical representation of the tube is obtained by giving the coordinates ω_{\min} and ω_{\max} of the points where λ_{\min} and λ_{\max} respectively intersect the ω -axis and the coordinates (ω_1, z_1) and (ω_2, z_2) where λ_{\min} and

TABLE 1. Coordinates $(\omega_{\min}, 0)$, $(\omega_{\max}, 0)$, (ω_1, z_1) and (ω_2, z_2) in kpc for a tube orbit with $Z = 10 \text{ kms}^{-1}$

$\Pi \setminus \Phi - \Phi_0$	kms^{-1}							
	-80	-60	-40	-20	0	20	40	
0	4.01, 0.00	4.92, 0.00	5.91, 0.00	7.03, 0.00	8.32, 0.00	8.50, 0.00	8.50, 0.00	
	8.50, 0.00	8.50, 0.00	8.50, 0.00	8.50, 0.00	8.50, 0.00	9.86, 0.00	11.77, 0.00	
	4.01, 0.22	4.91, 0.25	5.90, 0.28	7.02, 0.30	8.31, 0.33	8.49, 0.32	8.49, 0.30	
	8.49, 0.46	8.49, 0.43	8.49, 0.39	8.49, 0.37	8.49, 0.34	9.85, 0.37	11.77, 0.41	
30	3.97, 0.00	4.85, 0.00	5.80, 0.00	6.79, 0.00	7.69, 0.00	8.15, 0.00	8.31, 0.00	
	8.63, 0.00	8.66, 0.00	8.71, 0.00	8.85, 0.00	9.26, 0.00	10.36, 0.00	12.15, 0.00	
	3.97, 0.22	4.85, 0.25	5.79, 0.27	6.78, 0.29	7.68, 0.31	8.14, 0.31	8.30, 0.29	
	8.61, 0.47	8.65, 0.44	8.71, 0.40	8.84, 0.38	9.25, 0.37	10.35, 0.39	12.14, 0.42	
60	3.86, 0.00	4.68, 0.00	5.53, 0.00	6.35, 0.00	7.07, 0.00	7.58, 0.00	7.88, 0.00	
	9.00, 0.00	9.11, 0.00	9.30, 0.00	9.63, 0.00	10.26, 0.00	11.39, 0.00	13.14, 0.00	
	3.86, 0.22	4.68, 0.24	5.52, 0.26	6.35, 0.27	7.07, 0.28	7.57, 0.28	7.87, 0.28	
	8.99, 0.49	9.10, 0.46	9.29, 0.43	9.62, 0.41	10.26, 0.41	11.38, 0.43	13.13, 0.46	

TABLE 2. Coordinates $(\bar{\omega}_{\min}, 0)$, $(\bar{\omega}_{\max}, 0)$, $(\bar{\omega}_1, z_1)$ and $(\bar{\omega}_2, z_2)$ in *kpc* for a tube orbit with $Z = 50 \text{ kms}^{-1}$

$\Pi \setminus \bar{\phi} - \bar{\phi}_0$ <i>kms</i> ⁻¹	-80	-60	-40	-20	0	20	40
0	4.36, 0.00	5.26, 0.00	6.25, 0.00	7.38, 0.00	8.50, 0.00	8.50, 0.00	8.50, 0.00
	8.50, 0.00	8.50, 0.00	8.50, 0.00	8.50, 0.00	8.70, 0.00	10.29, 0.00	12.28, 0.00
	4.19, 1.26	5.08, 1.37	6.08, 1.49	7.21, 1.61	8.33, 1.72	8.35, 1.60	8.37, 1.49
	8.15, 2.42	8.21, 2.20	8.26, 2.01	8.30, 1.85	8.52, 1.76	10.11, 1.93	12.09, 2.15
30	4.32, 0.00	5.18, 0.00	6.11, 0.00	7.08, 0.00	7.86, 0.00	8.20, 0.00	8.33, 0.00
	8.64, 0.00	8.67, 0.00	8.75, 0.00	8.92, 0.00	9.47, 0.00	10.74, 0.00	12.64, 0.00
	4.14, 1.25	5.01, 1.35	5.94, 1.45	6.91, 1.55	7.70, 1.59	8.06, 1.54	8.20, 1.46
	8.29, 2.46	8.38, 2.24	8.50, 2.07	8.71, 1.94	9.28, 1.91	10.55, 2.02	12.45, 2.22
60	4.19, 0.00	4.98, 0.00	5.80, 0.00	6.58, 0.00	7.23, 0.00	7.67, 0.00	7.93, 0.00
	9.04, 0.00	9.17, 0.00	9.38, 0.00	9.78, 0.00	10.51, 0.00	11.76, 0.00	13.64, 0.00
	4.02, 1.21	4.81, 1.30	5.63, 1.38	6.42, 1.44	7.08, 1.46	7.53, 1.44	7.81, 1.39
	8.67, 2.57	8.86, 2.37	9.12, 2.22	9.54, 2.13	10.30, 2.12	11.55, 2.21	13.43, 2.39

respectively intersect the sheet ν_{\max} in the region $z \geq 0$. They are given or peculiar velocity components $|\Pi|$, $\bar{\phi} - \bar{\phi}_0$, adopted $\bar{\phi}_0 = 220 \text{ kms}^{-1}$, at a position $\bar{\omega} = 8.5 \text{ kpc}$ on the B-axis for $Z = 10 \text{ kms}^{-1}$ and $Z = 50 \text{ kms}^{-1}$ in tables 1 and 2 respectively.

II. THE ORBIT OF THE SUN

A more detailed study has been performed of the orbit of the Sun around the galactic rotation axis. A position $\bar{\omega}_\odot = 8.5 \text{ kpc}$, $z_\odot = 0.007 \text{ kpc}$

and peculiar velocity components $\Pi = -9 \text{ kms}^{-1}$, $\bar{\phi} - \bar{\phi}_0 = 12 \text{ kms}^{-1}$, $Z = 7 \text{ kms}^{-1}$ (Delhaye 1965) have been adopted for the Sun. Table 3 gives the values of the first integrals I_1 , I_2 and I_3 , the limiting coordinate surfaces λ_{\min} , λ_{\max} and

TABLE 3. Values of the first integrals and the limiting coordinates of the orbit of the Sun

$$I_1 = -141365.19 \text{ km}^2 \text{ s}^{-2}$$

$$I_2 = 1972.00 \text{ kms}^{-1} \text{ kpc}$$

$$I_3 = 3590.76 \text{ km}^2 \text{ s}^{-2} \text{ kpc}^2$$

$$\lambda_{\min} = 71.796 \text{ kpc}^2$$

$$\lambda_{\max} = 86.903 \text{ kpc}^2$$

$$\nu_{\max} = 0.000557 \text{ kpc}^2$$

$$(\bar{\omega}_{\min}, 0) = (8.43, 0.00) \text{ kpc}$$

$$(\bar{\omega}_{\max}, 0) = (9.28, 0.00) \text{ kpc}$$

$$(\bar{\omega}_1, z_1) = (8.42, 0.23)$$

$$(\bar{\omega}_2, z_2) = (9.28, 0.25)$$

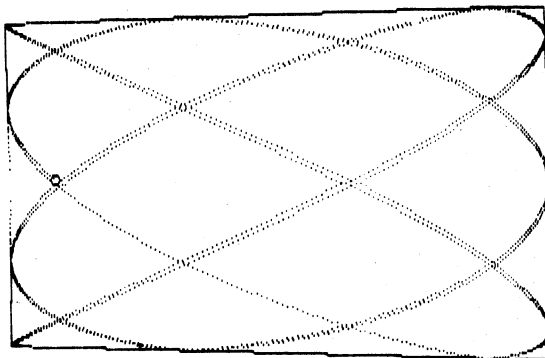


FIGURE 1. Projection of the orbit of the Sun on a meridional plane

TABLE 4. Position and velocity of the Sun along its orbit

t 10 ⁶ years	ϖ kpc	ϕ degrees	z kpc	Π	$\bar{\Phi}$ kms ⁻¹	Z	λ_{\min}	λ_{\max}	$z=0$	v_{\max}	revol.
0.00	8.50	0.00	0.007	-9.0	232.00	7.00					
14.71	8.43	23.63	0.11	0.3	234.01	6.25	*				
45.43	8.70	72.52	0.23	14.7	226.79	-0.45				*	
97.63	9.28	146.15	-0.00	0.5	212.49	-6.43			*		
98.20	9.28	146.92	-0.00	-0.2	212.49	-6.43		*			
150.34	8.71	220.30	-0.23	-14.9	226.38	1.23				*	
181.72	8.43	270.17	-0.10	0.4	234.01	6.25	*				
196.46	8.50	294.02	0.00	8.9	232.04	7.01			*		
240.50	9.12	360.00	0.24	11.8	216.33	2.15					*
246.60	9.18	368.29	0.25	9.5	214.78	-0.54				*	
265.21	9.28	393.44	0.20	-0.1	212.54	-3.71		*			
300.42	8.98	441.82	-0.00	-14.7	219.63	-6.64			*		
348.35	8.42	516.07	-0.23	-0.6	234.07	0.88				*	
348.74	8.42	516.70	-0.23	0.4	234.08	0.94	*				
395.76	8.95	589.77	0.00	15.0	220.25	6.66			*		
432.21	9.28	639.94	0.21	-0.2	212.54	3.60		*			
449.98	9.19	663.95	0.25	-9.0	214.52	-1.03				*	
488.60	8.65	720.02	0.08	-13.9	228.00	-6.60					*
499.86	8.51	737.55	-0.00	-9.6	231.16	-7.00			*		
515.73	8.43	763.20	-0.11	0.3	234.01	-6.24	*				
546.32	8.69	811.89	-0.23	14.7	226.84	0.45				*	
598.48	9.28	885.49	0.00	0.6	212.49	6.43			*		
599.21	9.28	886.46	0.00	-0.2	212.49	6.43		*			
651.27	8.71	959.74	0.23	-14.9	226.38	-1.23				*	
682.73	8.43	1009.72	0.10	0.4	234.01	-6.30	*				
697.36	8.50	1033.39	-0.00	8.9	232.06	-7.01			*		
727.80	8.92	1080.10	-0.19	15.1	221.00	-4.44					*

v_{\max} and the intersections $(\varpi_{\min}, 0)$, $(\varpi_{\max}, 0)$, (ϖ_1, z_1) and (ϖ_2, z_2) of the orbit of the Sun. Figure 1 shows the projection of the orbit of the Sun on a meridional plane, being this projection a quasi-periodic motion. Table 4 shows the changes of the position and the velocity of the Sun with time, indicating the consecutive crossings with the coordinate surfaces in its motion around the rotation axis of the Galaxy.

IV. SUMMARY

An Eddington type approximation of the galactic potential, with three first integrals in involution, has been used to determine the major properties of the orbits of stars in the Galaxy. The sizes of the tubes have been determined for different ranges of peculiar stellar velocities at the solar neighbourhood. Finally, a more detailed study of the orbit of the Sun has been performed.

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