

EXTENSIVE PHOTOMETRY OF V1838 AQL DURING THE 2013 SUPEROUTBURST

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ABSTRACT

We present an in-depth photometric study of the 2013 superoutburst of the recently discovered cataclysmic variable V1838 Aql and subsequent photometry near its quiescent state. A careful examination of the development of the superhumps is presented. Our best determination of the orbital period is $P_{\text{orb}} = 0.05698(9)$ days, based on the periodicity of early superhumps. Comparing the superhump periods at stages *A* and *B* with the early superhump value we derive a period excess of $\epsilon = 0.024(2)$ and a mass ratio of $q = 0.10(1)$. We suggest that V1838 Aql is approaching the orbital period minimum and thus has a low-mass star as a donor instead of a substellar object.

RESUMEN

Presentamos un estudio fotométrico detallado de la super-erupción de V1838 Aql, una variable cataclísmica recientemente descubierta, desde el máximo en 2013 hasta su regreso al mínimo. Examinamos en detalle la evolución de las superjorobas (*superhumps*). Determinamos el período orbital $P_{\text{orb}} = 0.05698(9)$ días a partir de la periodicidad de las superjorobas tempranas. Comparando los períodos de las superjorobas en las etapas *A* y *B* con el valor del período orbital, derivamos un valor del cambio en el período orbital de $\epsilon = 0.024(2)$ y un cociente de masa para el sistema de $q = 0.10(1)$. Sugerimos que V1838 Aql se está acercando al mínimo período orbital, por lo que la secundaria sería una estrella de baja masa y no un objeto sub-estelar.

Key Words: techniques: photometric — stars: novae, cataclysmic variables — stars: dwarf novae — stars: individual: V1838 Aql

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1. INTRODUCTION

Cataclysmic variables (CVs) are close binary systems in which a white dwarf (WD) accretes from a low-mass star via Roche-lobe overflow, often creating an accretion disc (for a review see Warner 1995). A large fraction of CVs belong to the subclass of dwarf novae (DNe). They undergo recurrent outbursts with typical amplitudes of $\approx 2 - 6$ mag in the optical, which are commonly accepted to be caused by a thermal-viscous instability in the disc (Osaki 1974). In addition, SU UMa-type DNe (the subclass of DNe systems with short orbital periods, $P_{\text{orb}} < 2.5$ hr) exhibit occasional eruptions that are less frequent, longer lasting, and slightly brighter (by $\approx 0.5 - 1.0$ mag) than the normal out-

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bursts. The key feature during these so-called superoutbursts is the presence of superhumps a modulation in the light curve with an underlying periodicity, P_{sh} , a few percent longer than the orbital period. They are thought to arise from a precessing non-axisymmetric disc (Vogt 1982), with the eccentricity being produced by the tidal instability developed at the radius of the 3:1 resonance (Whitehurst 1988). The analysis of the timing and evolution of such light oscillations provides estimates of the system’s parameters using empirical (Patterson et al. 2005) and theoretical (Kato & Osaki 2013) relationships between the superhump/orbital period excess, $\epsilon \equiv (P_{\text{sh}} - P_{\text{orb}})/P_{\text{orb}}$, and the binary mass ratio $q \equiv M_2/M_1$. Several $\epsilon(q)$ relations have been proposed by these authors based on different stages of the superhumps, although Otulakowska-Hypka et al. (2016) show that the scatter in the ϵ - q diagram is considerable for very short orbital periods. This is due to a lack of objects with dynamically confirmed small values of q (Patterson 2011; Kato & Osaki 2013).

Among the SU UMa-type systems, there is a large group that accumulates around the minimum of the orbital period distribution of CVs ($P_{\text{orb}} \approx 78$ min) (Paczynski & Sienkiewicz 1981; Gänsicke et al. 2009; Knigge et al. 2011). These are binaries with extremely low mass-transfer rates, named WZ Sge-type objects, which are characterised by rare (commonly detected every ≈ 10 years), and large amplitude superoutbursts of duration of ≈ 30 days, caused by an instability in low viscosity accretion discs with $\alpha \approx 0.01 - 0.001$ (Smak 1993; Osaki 1994). Some of them are systems currently evolving towards longer periods and are collectively known as period bouncers (e.g. Littlefair, Dhillon & Martin 2003). These binaries are expected to harbour a substellar secondary companion (Howell et al. 1997), i.e. brown dwarfs (e.g. Littlefair et al. 2006; Harrison 2016; Hernández Santisteban et al. 2016; Neustroev et al. 2017).

Worth noting is that WZ Sge-type objects are characterised not only by a long superoutburst recurrence time in comparison to typical SU UMa stars, but also by the presence of early superhumps (double-wave modulation) during the first few days of the eruption, with a periodicity (P_{esh}) essentially equal to the orbital period of the binary (details in O’Donoghue et al. 1991; Kato 2015). The advent of all-sky surveys (e.g. Breedt et al. 2014) and worldwide citizen-telescope networks has contributed to the discovery of a large population of faint DNe. Among these discoveries, the elusive population of

short-period systems, in particular period bouncers, has been found and investigated (Patterson 2011; Coppejans et al. 2016; Otulakowska-Hypka et al. 2016).

The discovery of a new transient, initially proposed as possible nova, was reported by Itagaki on 2013 May 31. Henden¹⁶ pointed out that the colour indices of the object and the un-reddened field suggested a DN rather than a nova. As pointed out by Hurst¹⁷, Kojima reported a pre-discovery image on 2013 May 30.721 UT, when the magnitude was at about 9.8 mag (un-filtered). A CBET report of this new DNe in Aquila, can be found in Itagaki (2013).

Although a preliminary analysis of the behaviour of V1838 Aql (originally designated as PNV J19150199+0719471) was published by Kato et al. (2014), we present here a full analysis of the superhump behaviour based on our extensive data.

In § 2 we present the observations and their reduction methods. The photometric data and the period analysis are presented in § 3, while in § 4 we address the discussion of our results. We present our conclusions in § 5.

2. OBSERVATIONS AND REDUCTION

Photometric observations in the V band were obtained in 2013 during the nights of June 3, 4, 5, 6, 17, 18, 28 and September 2 and 25 at the 0.84 m telescope of the Observatorio Astronómico Nacional at San Pedro Mártir (SPM). We used the Blue-ESOPO CCD detector¹⁸ on a 2×2 binning configuration. The exposure time of the SPM observations varied between 10 and 30 s. In addition, time-series photometry of the superoutburst was obtained from 10 observatories of the Center for Backyard Astrophysics (CBA) – a network of small (0.2 – 0.4 m) telescopes that covers a wide range in terrestrial longitude. Skillman & Patterson (1993) and de Miguel et al. (2016) describe the methods and observing stations of the CBA network. These observations amounted to 162 separate time-series during 58 nights from June 1 to August 2, 2013, and the typical exposure time ranged from 20 to 120 seconds, depending on the brightness of V1838 Aql. Nearly half of these observations were obtained in V light, while the rest (mainly during the post-outburst regime) were unfiltered. We did not attempt any absolute calibration of the data during the eruption, but the magnitude scale is expected to resemble

¹⁶<http://ooruri.kusastro.kyoto-u.ac.jp/mailarchive/vsnet-alert/15772>.

¹⁷www.theastronomer.org/tacirc/2013/e2919.txt.

¹⁸<http://www.astrossp.unam.mx/indexspm.html>.

closely V magnitudes with a zero-point uncertainty of ≈ 0.05 mag. Further observations in the R band using the 2.1 m telescope at SPM during 2018, July 18 were conducted. Unfortunately, the weather was unstable and we only managed to obtain differential photometry over three orbital cycles. In the following, we report times and refer to specific dates in a truncated form defined as $\text{HJD} - 2,456,000$.

3. PHOTOMETRY AND PERIOD ANALYSIS

3.1. Photometric Observations

Most of our photometric observations come from the CBA network, with additional V -band observations obtained with the 0.84 m telescope at SPM at some critical stages of the outburst and during the late decline.¹⁹ We observed the typical pattern seen in SU UMa stars during superoutburst: a plateau phase –lasting ≈ 25 days, from HJD 444 to 469– where the mean brightness varied smoothly from 10.5 to 13.0 mag, followed by a rapid decline (≈ 3 mag in 2 days) at the end of the main eruption. The subsequent fading towards quiescence occurred at a rather low rate (≈ 0.035 mag d^{-1}), and even 3.5 years after the end of the eruption, the system was found to be ≈ 0.5 mag above the pre-eruption quiescent brightness. However, based on the observations obtained with the NTT telescope (La Palma, Spain), we confirmed that the object had reached the pre-outburst level by June 2017. These observations and the general spectral distribution at quiescence have no further relevance here and will be discussed in a future publication.

3.2. Full Analysis of the Different Stages of the Superhump

Our primary tool for studying periodic signals was the Period 04 package (Lenz & Breger 2005). First, we subtracted the mean and (linear) trend from each individual light curve and formed nightly-spliced light curves. Then, after combining light curves from adjacent nights, a search for periodic signals was done. This approach allows us to improve the frequency resolution, but it has to be implemented with caution, since variations in the amplitude and/or the period of the modulation –both effects known to afflict erupting DNe– can distort the outcome of the frequency analysis.

A general view of the superhump transitions can be looked up by identifying the different stages of the

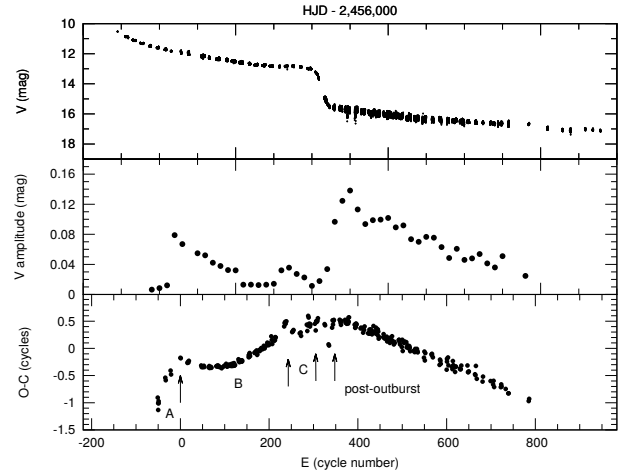


Fig. 1. Photometric behaviour of V1838 Aql during its superoutburst in 2013. Top frame: global light curve of the observational campaign. Middle frame: amplitude variations of the superhumps along the superoutburst (see text). Bottom frame: O–C diagram for the superhump maxima in the HJD 449.6–498.6 day interval with respect to the ephemeris given in equation (1). The arrows indicate the (approximate) location of transitions between different regimes of superhump period variations (see text).

superhumps: early superhumps, Stages *A*, *B* and *C* as well as the post-outburst stage (see Kato et al. 2009, 2014, for this terminology in our general discussion). Thus, we looked at the time variations of the superhump period and its amplitude by examining the variation in time with respect to a well-defined feature of the superhump signal.

After this general analysis, a detailed explanation of these stages was made. First, we derived the timings of superhump maxima. A total of 310 times of superhump maxima was identified in the light curves in the interval HJD 449.6–498.6 days. These maxima are shown in Table 1. The early stage of the eruption was not considered, since the signal there was of very low amplitude and individual maxima were not well defined. A linear regression to these timings provides the following test ephemeris:

$$T_{\text{max}}(\text{HJD}) = 2,456,452.8035(23) + 0.058191(6) E. \quad (1)$$

The top panel of Figure 1 displays the general photometric behaviour of the system during our campaign. Next, on the middle panel of Figure 1 is presented the variation of the amplitude of the superhump modulation, defined as the semi-amplitude of the sine wave that best fits the nightly photometric data.

¹⁹All CBA and SPM data are available on request.

TABLE 1
TIMES OF SUPERHUMP MAXIMA OF V1838 Aql DURING THE 2013 SUPEROUTBURST.^a

E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d
-51	449.7831	-0.906	116	459.5372	-0.283	189	463.8032	0.028	345	472.9088	0.506
-50	449.8282	-1.131	116	459.5365	-0.295	190	463.8598	-0.001	346	472.9681	0.524
-50	449.8354	-1.007	117	459.5944	-0.301	190	463.8607	0.016	359	473.7247	0.527
-49	449.8934	-1.011	117	459.5939	-0.309	195	464.1540	0.056	360	473.7820	0.511
-49	449.8956	-0.972	118	459.6513	-0.322	196	464.2150	0.104	361	473.8381	0.474
-34	450.7936	-0.541	120	459.7703	-0.278	197	464.2728	0.097	362	473.8938	0.431
-33	450.8491	-0.586	121	459.8289	-0.271	198	464.3288	0.060	363	473.9561	0.502
-22	451.4995	-0.410	122	459.8842	-0.320	199	464.3889	0.092	371	474.4228	0.523
-21	451.5535	-0.481	123	459.9432	-0.307	200	464.4455	0.065	372	474.4782	0.475
0	452.7934	-0.174	131	460.4149	-0.200	205	464.7456	0.222	373	474.5387	0.514
1	452.8512	-0.181	132	460.4708	-0.239	206	464.7997	0.151	374	474.5940	0.464
16	453.7190	-0.268	132	460.4718	-0.223	207	464.8537	0.080	375	474.6515	0.453
17	453.7782	-0.250	133	460.5281	-0.254	212	465.1482	0.141	376	474.7119	0.490
18	453.8358	-0.261	133	460.5286	-0.247	213	465.2072	0.154	377	474.7702	0.492
19	453.8959	-0.228	134	460.5883	-0.220	214	465.2684	0.205	377	474.7705	0.497
47	455.5203	-0.313	134	460.5883	-0.220	215	465.3263	0.201	377	474.7706	0.500
48	455.5766	-0.345	135	460.6459	-0.231	234	466.4437	0.403	378	474.8263	0.456
49	455.6356	-0.331	136	460.7049	-0.218	235	466.5055	0.465	378	474.8280	0.486
50	455.6938	-0.331	137	460.7625	-0.227	236	466.5650	0.488	379	474.8868	0.497
51	455.7512	-0.345	138	460.8201	-0.237	237	466.6212	0.453	379	474.8862	0.486
52	455.8099	-0.336	139	460.8780	-0.243	239	466.7401	0.497	379	474.8909	0.566
53	455.8690	-0.321	139	460.8787	-0.230	252	467.4850	0.298	380	474.9493	0.570
54	455.9259	-0.342	139	460.8792	-0.221	253	467.5438	0.309	393	475.6979	0.435
64	456.5073	-0.351	154	461.7554	-0.164	254	467.6029	0.324	394	475.7561	0.435
65	456.5671	-0.323	154	461.7576	-0.126	255	467.6619	0.338	394	475.7527	0.376
66	456.6232	-0.361	155	461.8140	-0.158	269	468.4730	0.277	395	475.8151	0.449
69	456.7989	-0.341	156	461.8764	-0.085	271	468.5903	0.292	395	475.8118	0.392
70	456.8561	-0.358	156	461.8756	-0.098	272	468.6448	0.229	395	475.8140	0.430
81	457.4979	-0.328	157	461.9311	-0.144	285	469.4134	0.436	396	475.8725	0.436
82	457.5560	-0.330	167	462.5166	-0.083	286	469.4648	0.320	396	475.8737	0.456
83	457.6146	-0.323	168	462.5755	-0.070	287	469.5288	0.420	397	475.9309	0.439
84	457.6733	-0.314	169	462.6334	-0.076	288	469.5973	0.598	406	476.4504	0.366
85	457.7303	-0.335	170	462.6916	-0.075	289	469.6537	0.567	407	476.5093	0.380
86	457.7893	-0.320	171	462.7489	-0.090	290	469.7051	0.451	411	476.7418	0.374
87	457.8469	-0.332	171	462.7500	-0.072	291	469.7620	0.428	412	476.7950	0.288
88	457.9044	-0.343	171	462.7514	-0.047	292	469.8186	0.400	413	476.8565	0.345
92	458.1365	-0.354	172	462.8101	-0.039	292	469.8215	0.450	414	476.9143	0.338
93	458.1941	-0.365	172	462.8074	-0.086	293	469.8789	0.437	415	476.9763	0.404
97	458.4284	-0.338	172	462.8078	-0.078	303	470.4617	0.452	423	477.4346	0.280
99	458.5459	-0.318	173	462.8644	-0.106	305	470.5709	0.329	425	477.5504	0.270
102	458.7208	-0.313	173	462.8635	-0.122	306	470.6391	0.501	426	477.6130	0.346
103	458.7807	-0.283	178	463.1596	-0.033	308	470.7548	0.489	427	477.6698	0.321
104	458.8372	-0.313	179	463.2169	-0.049	309	470.8143	0.510	429	477.7832	0.270
105	458.8942	-0.333	180	463.2767	-0.021	310	470.8750	0.554	429	477.7877	0.348
106	458.9522	-0.336	181	463.3344	-0.030	325	471.7415	0.446	430	477.8417	0.275
110	459.1872	-0.298	182	463.3946	0.005	326	471.7985	0.424	430	477.8434	0.305
111	459.2441	-0.321	184	463.5105	-0.003	334	472.2435	0.072	431	477.9022	0.315
112	459.3025	-0.316	185	463.5680	-0.015	335	472.3005	0.051	432	477.9570	0.257
113	459.3617	-0.300	186	463.6273	0.003	342	472.7271	0.381	435	478.1406	0.412
115	459.4778	-0.305	187	463.6863	0.019	343	472.7931	0.516	436	478.1927	0.308
115	459.4767	-0.323	188	463.7462	0.047	344	472.8467	0.437	437	478.2500	0.293

TABLE 1. CONTINUED

E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d	E ^b	T _{max} ^c	O-C ^d
438	478.3092	0.310	482	480.8586	0.120	535	483.9315	-0.073	631	489.4996	-0.386
440	478.4235	0.273	483	480.9154	0.097	549	484.7405	-0.170	632	489.5580	-0.382
441	478.4880	0.381	492	481.4423	0.151	550	484.8051	-0.060	633	489.6199	-0.318
442	478.5401	0.277	493	481.5009	0.159	551	484.8538	-0.222	647	490.4349	-0.313
443	478.6004	0.313	494	481.5564	0.112	552	484.9199	-0.087	648	490.4859	-0.437
444	478.6549	0.251	495	481.6122	0.071	566	485.7333	-0.109	649	490.5440	-0.438
445	478.7189	0.351	499	481.8423	0.026	567	485.7904	-0.128	650	490.6030	-0.424
446	478.7747	0.309	500	481.9082	0.158	567	485.7921	-0.098	664	491.4146	-0.476
447	478.8253	0.179	501	481.9544	-0.048	569	485.9045	-0.166	665	491.4719	-0.493
447	478.8274	0.215	504	482.1375	0.098	584	486.7694	-0.303	666	491.5400	-0.322
457	479.4087	0.204	505	482.1974	0.127	585	486.8333	-0.206	668	491.6425	-0.560
458	479.4702	0.262	509	482.4295	0.116	586	486.8901	-0.229	681	492.3960	-0.611
460	479.5823	0.188	510	482.4819	0.017	587	486.9451	-0.285	682	492.4595	-0.520
461	479.6400	0.180	511	482.5406	0.025	600	487.6953	-0.392	699	493.4440	-0.602
464	479.8154	0.194	512	482.6006	0.056	601	487.7651	-0.193	705	493.7895	-0.665
465	479.8748	0.214	515	482.7719	-0.000	601	487.7622	-0.243	706	493.8516	-0.597
466	479.9279	0.127	516	482.8296	-0.008	603	487.8830	-0.167	707	493.9032	-0.711
470	480.1643	0.188	516	482.8326	0.042	604	487.9327	-0.312	722	494.7758	-0.715
471	480.2222	0.185	517	482.8892	0.015	612	488.3968	-0.337	723	494.8371	-0.662
472	480.2755	0.100	518	482.9433	-0.054	613	488.4523	-0.383	724	494.8931	-0.700
474	480.3975	0.197	526	483.4092	-0.047	614	488.5193	-0.231	733	495.4140	-0.747
475	480.4534	0.157	527	483.4678	-0.040	615	488.5730	-0.308	739	495.7585	-0.828
476	480.5116	0.157	528	483.5244	-0.068	618	488.7454	-0.347	740	495.8167	-0.828
477	480.5736	0.224	529	483.5860	-0.010	619	488.7985	-0.434	785	498.4271	-0.968
478	480.6261	0.125	532	483.7593	-0.032	620	488.8621	-0.341	786	498.4876	-0.929
480	480.7420	0.117	533	483.8188	-0.010	621	488.9207	-0.333			
481	480.7976	0.072	534	483.8748	-0.047	630	489.4422	-0.372			

^aIndividual errors in these timings are not explicitly included here.

^bE (cycle number).

^cSuperhump maxima expressed as HJD - 2,456,000.

^dO - C value (in cycles) according to the ephemeris.

T_{max} (HJD) = 2,456,452.8035 + 0.058191 E.

The O–C residuals of the times of maximum light relative to the ephemeris given by equation 1 are shown in the lower panel of Figure 1. The resulting O–C diagram is complex, but it displays a number of features that are usually observed in other SU UMa-type systems (Kato 2015). Among the most relevant features visible in this diagram we point out the following:

1. During the first four days of the outburst a weak modulation (early superhumps) with a period $P_{\text{esh}} \approx 0.057$ d was visible in the light curve.
2. The onset of fully-grown (Stage *A* superhumps) took place in a short time-scale (≈ 2 d) and involved an increase in the amplitude of the modulations. Their (mean) period, $P_{\text{sh}(A)} \approx 0.059$ d, was longer than P_{esh} .
3. Once the superhump modulation reached full amplitude, the system entered Stage *B* where the amplitude of the superhump decreased slowly, and the mean period became shorter ($P_{\text{sh}(B)} \approx 0.058$ d). The upward curvature of the residuals during this stage (days HJD 449–466) signifies that the period of the superhumps was not constant, but increased over time. From a quadratic fit of the residuals in this interval, we find an increase rate of $dP_{\text{sh}(B)}/dt = 5.8(4) \times 10^{-5}$.
4. Before the end of the main eruption, the amplitude of the superhumps was found to grow larger (≈ 0.10 mag). The system entered Stage *C*, extending from day HJD 466 to the end of the main plateau (around day HJD 470), where the period of the superhump remained essentially constant ($P_{\text{sh}(C)} \approx 0.0582$ d).

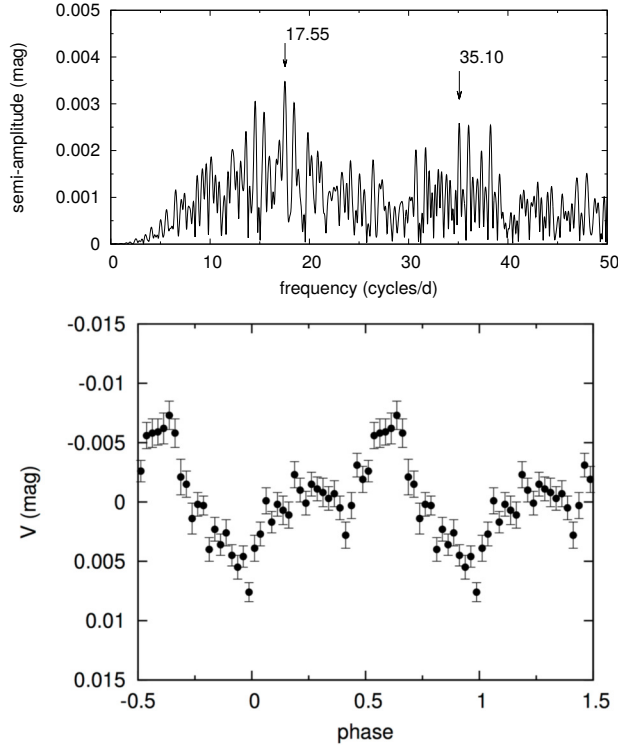


Fig. 2. Upper frame: Power spectrum during days HJD 445–448 (early superhumps), showing broad peaks centred at 17.55 cycles d^{-1} and its first harmonic. Lower frame: Waveform of the early superhumps obtained after folding the data with $P_{\text{esh}} = 0.05698$ d. The zero phase is arbitrary.

5. Worth noting is the increase of the amplitude variations as the system dropped by about 3 mag between Stage *C* and the post-outburst stage.
6. After the end of the main eruption (day \geq HJD 473) the superhumps were still visible with a significantly larger amplitude (≈ 0.2 mag), and with a period which remained constant for at least the subsequent ≈ 25 days. The period of the post-outburst modulation was shorter than P_{sh} .

A summary of the main periodicities along the eruption is given in Table 2, as found in the next subsections.

3.2.1. Early Superhumps

From the beginning of our campaign, a weak modulation of about 0.010 mag full amplitude was observed in the light curves. This signal persisted over days HJD 445–448. The power spectrum of the spliced light curve covering this 4-day segment

is shown in the upper frame of Figure 2. It is dominated by two broad peaks centred at frequencies 35.10(3) and 17.55(3) cycles d^{-1} . These signals were weak, with amplitudes of 0.0036 and 0.0030 mag, respectively. Although they were barely detected above the noise, we interpret them as a likely manifestation of early superhumps (Kato et al. 2014).

This photometric feature is known to be typical of WZ Sge-type stars, and is not shown by any other type of dwarf nova. Although its physical origin is still under debate, there is increasing observational evidence that its period (P_{esh}) is essentially equal to P_{orb} (Patterson et al. 1996; Kato 2015). A folded curve of the spliced light curve with $P_{\text{esh}} = 0.05698(9)$ d is also shown in the lower frame of Figure 2, which shows the double-humped pattern characteristic of early superhumps. The value of P_{esh} obtained in this paper is slightly different from, but consistent with, the value of 0.05706(2) d reported in Kato et al. (2014).

3.2.2. Common Superhumps

The double-humped pattern of the early superhumps turned into single-peaked humps on day HJD 449. Over days HJD 449–451, the mean amplitude was around ≈ 0.007 mag and the period was $\approx 4\%$ longer than the period found for early superhumps. The dominant signal occurred at 16.83(8) cycles d^{-1} corresponding to a period of 0.0594(3) d. Since the modulation was better defined in this 2-day interval, we were able to determine the times of maximum in the signal. We identified a total of 9 maxima in the HJD 449.6–451.6 day interval, and obtained a period of 0.05934(11) d (corresponding to a frequency of 16.85(3) cycles d^{-1}) from a linear regression. This value is fully consistent with the one found from the Fourier analysis. The modulation over this 2-day segment is interpreted as Stage-*A* superhumps. The period we find is close to, but slightly different from, the value of 0.05883(6) d reported in Kato et al. (2014).

Fully-grown, large-amplitude superhumps were finally observed on day HJD 452 (amplitude of 0.10 mag). As a representative example, we show in the upper frame of Figure 3 the light curve from day HJD 455. As the eruption proceeded, the mean amplitude of the superhumps decreased (as shown in the middle panel of Figure 3). The variation in amplitude was smooth in the HJD 452–466 day interval. We formed a spliced light curve in this interval, and obtained the power spectrum shown in the middle frame of Figure 3. The strongest signals occurred at $f_1 = 17.128(3)$ and $f_2 = 34.265(3)$ cycles d^{-1} .

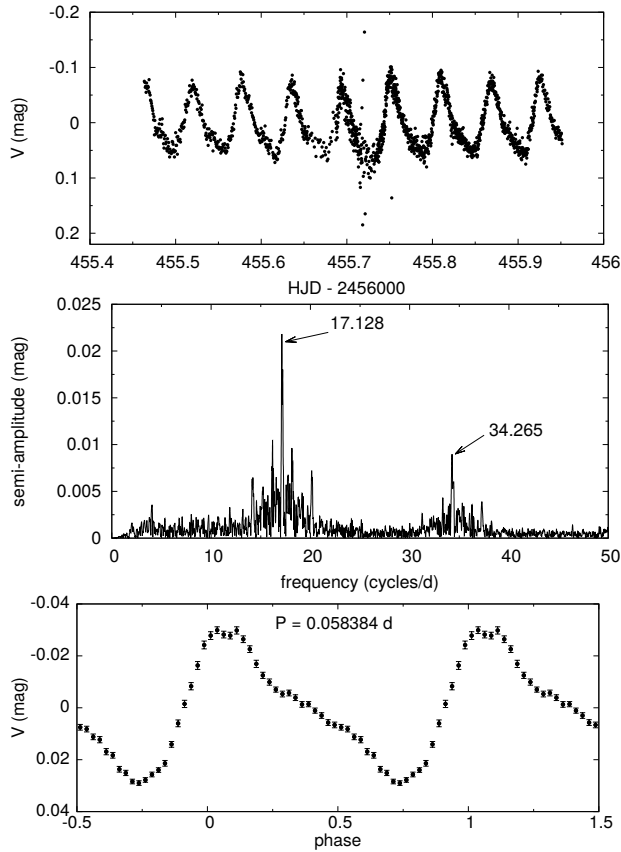


Fig. 3. Upper frame: A 12-hour spliced light curve obtained on day HJD 455, dominated by large-amplitude common superhumps. The zero level in the figure corresponds to $V \approx 12.1$ mag. Middle frame: Power spectrum during the common superhump era (days HJD 442–466), with main peaks centred at frequencies 17.128 and 34.265 cycles d^{-1} . Lower frame: Mean waveform of the common superhump obtained after folding the data on $P = 0.058384$ d. The zero phase is arbitrary.

They were interpreted as the frequency of Stage-*B* superhumps (period of $P_{\text{sh}} = 0.058384(10)$ d) and its first harmonic, respectively. Other (weaker) peaks, not shown in Figure 3, were found at $f_3 = 51.381(6)$ and $f_4 = 68.403(6)$ cycles d^{-1} . The mean waveform of the superhump modulation during this interval is shown in the lower panel of Figure 3.

We note that after subtracting the superhump signal and its harmonics, the power spectrum of the residual light curve showed peaks at 17.20 and 17.28 cycles d^{-1} . But we do not give any physical significance to these detections, and interpret them as the result of period and amplitude variations of the superhump wave during the eruption.

The amplitude of the superhumps increased around day HJD 466, and decreased thereafter until the end of the main eruption (Stage *C*). The

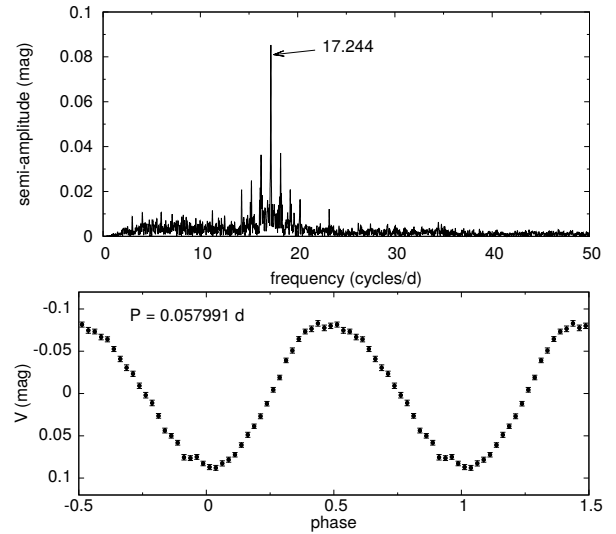


Fig. 4. Upper frame: Power spectrum after the main eruption (day HJD ≥ 473) showing a strong peak centred at 17.244(1) cycles d^{-1} (post-outburst superhump). Lower frame: Mean waveform of the post-outburst superhump, obtained after folding the data on $P = 0.057991$ d. The zero phase is arbitrary.

strongest signal in the power spectrum in the HJD 466–470 day interval occurred at 17.191(10) cycles d^{-1} , corresponding to a period of 0.05817(3) d, with additional peaks at higher harmonics.

3.2.3. Post-Outburst Stage

Once the main eruption was over, the light curve was still dominated by superhumps, but now with significantly larger amplitude (≈ 0.15 mag) which decreased slowly (≈ 0.013 mag d^{-1}). This behaviour remained essentially unchanged for nearly 20 days of our observations after the main fading.

We formed a spliced light curve including all the observations from day HJD ≥ 473 , and found a power spectrum (shown in the upper frame of Figure 4), with a peak at 17.244(1) cycles d^{-1} . This was interpreted as the frequency of the post-outburst superhump, and dominated the spectrum. Higher-order harmonics were also found, but their amplitude was very low (< 0.0065 mag). This signifies that the waveform of the post-outburst superhump was nearly sinusoidal. The lower frame of Figure 4 shows that this was indeed the case.

3.3. Photometry Near and at Minimum Light

As detailed in § 2, we took two runs near minimum light covering around one orbital period each. The light curves are shown in the upper panel of Figure 5. When folded with the orbital period

TABLE 2
MEAN PHOTOMETRIC PERIODS AND FREQUENCIES OF THE SUPERHUMP MODULATION.*

Time interval (HJD–2456000)	Period (d)	Frequency (cycles d ⁻¹)	Comments
445–448	0.05698(9)	17.55(3)	early superhumps
449–451	0.0594(3)	16.83(8)	Stage A
452–466	0.058384(10)	17.128(3)	Stage B
466–470	0.05817(3)	17.191(10)	Stage C
473–498	0.05799(1)	17.244(1)	post outburst
449–498	0.0581910(1)	17.1847(3)	mean superhump period

*Values correspond to the 2013 superoutburst of V1838 Aql. The errors in parentheses correspond to the last significant figures.

(0.05698 d), the light curves seem to be out of phase. But this is not surprising: over the 23 days (nearly 400 orbital cycles) elapsed between both runs, an uncertainty of 0.0001 days in P_{orb} involves an uncertainty of 0.7 in phase. We carried out a period analysis of both nights using the Phase Dispersion Minimization (PDM) technique (Stellingwerf 1978) in the *Peranso* package (Paunzen & Vanmunster 2016). This technique is frequently used to detect variations of superhumps in SU UMa systems (e.g. Kato et al. 2014). The lower frame in Figure 5 shows the results of combining the two nights with the best period estimate (0.0576 d) determined from the PDM technique. Assuming that the observed light comes from the accretion disc, the zero point obtained in this case is HJD 2456537.6946 (time of inferior conjunction of the secondary). The period found using the PDM method yielded a value which is still close to the post-outburst state, but the sinusoidal shape is gone. There was only a small peak around phase 0.25. No double modulation with orbital period was found as would be expected in a bounce-back object. Further observations in the *R* band were obtained on 2018, July 18 covering three orbital cycles. Since the night was not photometric, we were unable to make absolute calibrations and only differential photometry is shown in Figure 6. No obvious orbital modulation was detected within the individual errors, which are rather large (≈ 0.03 mag).

4. DISCUSSION

The values derived for the superhump period in Stages A and B allowed us to estimate the mass ratio of the system through the known “superhump excess” $\epsilon - q$ relations (Patterson 1998, 2011; Kato & Osaki 2013). Considering the lack of a reliable determination of the orbital period from spectroscopic obser-

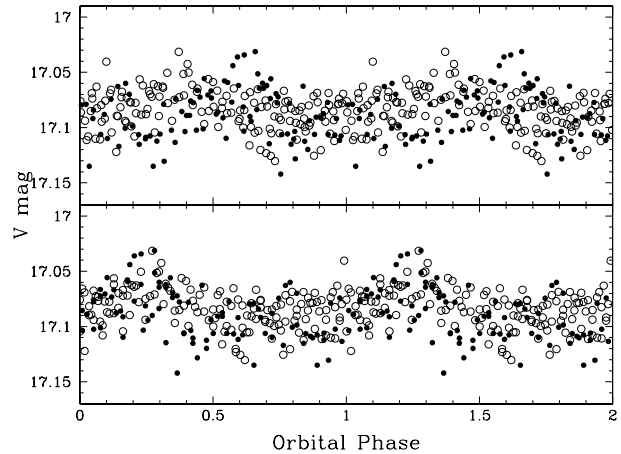


Fig. 5. Upper panel: *V* light curves obtained near minimum light in 2013. Open dots are from September 2 and filled dots from September 25. The size of the points corresponds to the mean individual errors. The orbital phases have been taken from the ephemeris obtained in this paper. Lower panel: The two nights folded with the ephemeris obtained by using the PDM technique (see more details in text).

vations, we assume here that P_{orb} is equal to the period of early superhumps. We find $\epsilon_A = 0.042(5)$ and $\epsilon_B = 0.024(2)$. Thus, our estimates for the mass ratio are $q_A = 0.12(2)$ and $q_B = 0.10(1)$, respectively. Comparing these two values with the $\epsilon - q$ relation shown in Bakowska et al. (2017, Figure 19) we can see clearly that ϵ_B is well within the expected value, while ϵ_A is not. This is further supported by using the updated Stolz & Schoembs (1984) relation in Otulakowska-Hypka et al. (2016, equation 4), which for our assumed orbital period gives $\epsilon = 0.019(10)$. Although we are inclined to use the Stage B results, we point out that both values are suggestive of a

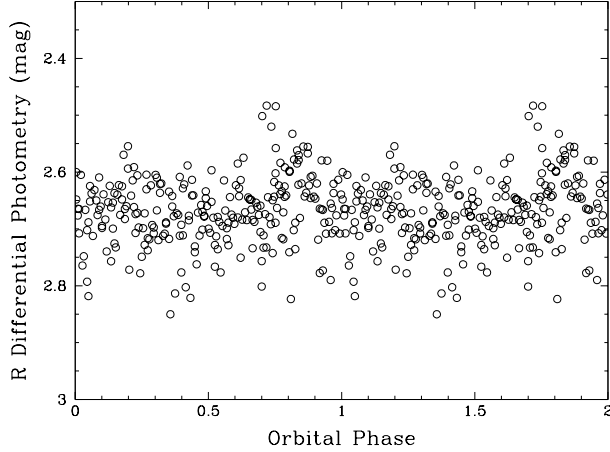


Fig. 6. Differential R light curve obtained at minimum light in July 18, 2018. The size of the individual errors (≈ 0.03 mag) is rather large. The orbital phases have been taken from the ephemeris obtained in this paper. No obvious orbital modulation, within the errors, is detected at this level.

low-mass donor, although as pointed out in § 1 (e.g. Otulakowska-Hypka et al. 2016, see their Figure 7), empirical relations ϵ - q at low q values may carry large systematic uncertainties. For a typical white dwarf with mass $\approx 0.8M_{\odot}$ (Zorotovic, Schreiber & Gänsicke 2011) and mass ratio $q \approx 0.1$, the mass of the secondary is very close to the substellar limit i.e. $0.072M_{\odot}$ (Chabrier & Baraffe 2000). Further characterisation of systems like V1838 Aql will allow us to discern empirically where this limit lies for mass-losing donors.

It has been noted by many authors, both theoretically and observationally, that the CV orbital period distribution should present a sharp cut-off at about ≈ 80 min, usually termed as the minimum period (e.g. Rappaport et al. 1982; Ritter & Kolb 1998; Gänsicke et al. 2009). V1838 Aql has an orbital period of about 82 min, very close to the minimum period, which makes it difficult to discern whether it is approaching to, or receding from, this minimum orbital period. Before asserting its true nature, we could look at some observational features in those CVs systems around the minimum orbital period. Most of these systems possess WZ Sge-like features. Their optical spectra are mostly dominated by the white dwarf and accretion disc itself, with no visible features from the donor. Since the rate of accretion is an order of magnitude smaller than that for systems before reaching the minimum period ($\dot{m} \approx 10^{-11}M_{\odot} \text{ yr}^{-1}$), the accretion discs become very faint, and the broad absorption lines of the white dwarf become visible below $\approx 5000 \text{ \AA}$ e.g.

WZ Sge (Howell et al. 2008). However, the donor’s observed properties should vary significantly for systems with the same orbital period but evolving towards or away from the period minimum. This is a consequence of the donor’s temperature steep relationship as a function of orbital period (e.g. Knigge et al. 2011). From the superhump analysis presented here, we suspect that the system is approaching the period minimum and thus, the donor is probably a late-M dwarf with an observed effective temperature of $\approx 2400 \text{ K}$ (Knigge et al. 2011).

The donor of V1838 Aql is therefore an ideal candidate for NIR time-resolved spectroscopy (e.g. SDSS J143317.78+101123.3, Hernández Santisteban et al. 2016), which would render a fully independent measurement of the orbital period and the mass ratio, to confirm or reject its substellar nature. This is particularly important since few low- q systems have been observed in outburst and for which a dynamical measurement of their components is feasible. (Figure 3 in Kato & Osaki 2013). Thus, V1838 Aql could be a system to calibrate the empirical superhump relations in this poorly-explored region of parameter space.

5. CONCLUSIONS

We have presented a long-term study of the 2013 superoutburst of V1838 Aql from its peak to quiescence. Our main results are as follows:

- The observed early superhumps suggest an orbital period of $P_{\text{orb}} = 0.05698(9)$ d, which locates V1838 Aql close to the minimum of the orbital period distribution in CVs.
- From Stages *A* and *B* and early superhump periods, we found the mass ratio to be $q_A = 0.12(2)$ and $q_B = 0.10(1)$, respectively.
- Based on the obtained values of the mass ratio, we claim that the donor in V1838 Aql is a low-mass star rather than a substellar object, and that the system is approaching the period minimum.

Given the long interval between outbursts in low q systems, it is of paramount importance to confirm by dynamical methods the orbital parameters of such systems. This would indicate which systems may be used in order to calibrate the empirical superhump excess relations.

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