THE ADF AND THE t^2 FORMALISM IN H II REGIONS BASED ON THE UPPER MASS LIMIT OF THE IMF FOR THE MW

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ABSTRACT

We study in depth the abundance discrepancy problem in H II regions, this time from a different perspective than the usual one: by studying the effect of the upper mass limit (M_{up}) of the initial mass function (IMF) on the O, C, and He predicted by chemical evolution models for the Milky Way. We use abundances determined with the direct method (DM) and with the temperature independent method (TIM). We compare the predicted abundances at the present time with observations of Orion, M17, and M8 to determine the M_{up} value of the galactic IMF. From the DM abundances, the models predict an $M_{up} = 25 - 45 M_{\odot}$, while from the TIM, CEMs derive an $M_{up} = 70 - 110 M_{\odot}$. Spiral galaxies with the stellar mass and star formation rate of the MW are predicted to have an $M_{up} \approx 100 M_{\odot}$. These results support that abundances derived from the TIM are better than those derived from the DM.

RESUMEN

Estudiamos el problema de la discrepancia de abundancias en regiones H II a partir del efecto que tiene la masa superior (M_{up}) de la función inicial de masa (IMF) en la producción de O, C y He obtenida por modelos de evolución química de nuestra galaxia. Comparamos las abundancias determinadas por el método directo (DM) y por el método independiente de la temperatura (TIM) en Orión, M17 y M8 con las abundancias modeladas actuales para determinar el valor de la M_{up} de la IMF Galáctica. Al usar las abundancias a partir del DM se requiere que $25 < M_{up} < 40 M_{\odot}$, mientras que con las abundancias obtenidas a partir del TIM se requiere que $70 < M_{up} < 110 M_{\odot}$. Para galaxias espirales con masa estelar y tasa de formación estelar similar a la VL se ha predicho una $M_{up} \approx 100 M_{\odot}$. Este resultado implica que las abundancias obtenidas por el TIM son más adecuadas que las obtenidas por el DM.

Key Words: Galaxy: evolution — HII regions — ISM: abundances

1. INTRODUCTION

The study of the chemical composition of H II regions is crucial for the understanding of the chemical composition of the universe; it can provide observational constraints required by models of galactic chemical evolution. The proper determination of the oxygen abundance in H II regions is critical to be able to compare our determinations with other branches of astrophysics, as well as to be able to present a coherent model of the evolution of the universe.

Although a comprehensive chemical composition is desired, frequently studies are only able to determine a few chemical elements; in fact many works focus only on the oxygen abundance. A single well determined element can be quite useful, since most elements are expected to behave in an orchestrated fashion, and oxygen is the ideal candidate, since it is expected to comprise about half of the heavy element abundance; not only that, but is the easiest element to study since it is the only element that produces bright optical lines for each of its main ionic species.

There are three competing methods to derive ionic O^+ and O^{++} abundances, relative to the H⁺ abundance, in H II regions: (i) The use of [O II] and [O III] collisionally excited lines (CELs) together with the H Balmer lines; these lines are relatively

easy to observe. Unfortunately their emmisivities depend strongly on the local temperature (in fact, this specific characteristic of CELs is used to determine the characteristic temperature of photoionized regions); this is known as the direct method (DM). (ii) The use of OI and OII recombination lines (RLs) together with the H Balmer lines (which are also RLs); this strategy has the advantage that the ratios of any pair of RLs are almost independent of the temperature structure; however, RLs of heavy elements are quite faint and harder to work with. (iii) The use of CELs with the t^2 formalism introduced by Peimbert (1967): Peimbert & Costero (1969) in which the effect of the temperature structure on the emission lines is taken into account. Recent reviews of these methods have been presented by Pérez-Montero (2017); Peimbert et al. (2017); Peimbert (2019); García-Rojas (2020).

Abundances derived from either RLs or the t^2 formalism usually agree with each other and are usually 0.2 to 0.3 dex higher than those determined using the DM (e.g. Esteban et al. 2009; Peimbert et al. 2017; Peimbert 2019; Carigi et al. 2019a). We have defined the temperature independent method (TIM) as abundances derived from either RLs, the t^2 formalism, or their average (Carigi et al. 2019a). On the other hand, the ratio between RL abundances and DM abundances (TIM abundances to DM abundances) is called the abundance discrepancy factor (ADF): $ADF(X^{+i}) = X_{\rm RLs}^{+i}/X_{\rm CELs}^{+i}$ (Tsamis et al. 2003).

Whether the origin of the abundance discrepancy is (only) due to temperature fluctuations (Peimbert 1967) or to chemical inhomegeneities (first proposed by Torres-Peimbert et al. 1990) is still under discussion (see e.g., Esteban, Toribio San Cipriano & García-Rojas 2018; García-Rojas et al. 2019). Other hypotheses have not been very successful and have been proposed over the years, such as the κ distribution of electrons (Nicholls, Dopita & Sutherland 2012), or uncertainties in the atomic data (e.g., Rodríguez & García-Rojas 2010), but have been discarded (García-Rojas et al. 2019). Anyway, it is beyond the scope of this paper to discuss the origin of this long-standing problem.

Oxygen observations of H II regions are limited to the O⁺ and O⁺⁺ gaseous components. In general, it is not necessary to correct for other ionization stages where O⁺³ is limited to $\approx 1-2\%$ and O⁰ is considered to be outside the relevant volume of the H II region; and, when comparing abundances between several H II regions, this is enough. However, when comparing with other types of objects, or when trying to model the evolution of a galaxy, it is of critical importance to correct for oxygen atoms trapped in dust grains, which are estimated to be $\approx 25\%$ of the total oxygen in H II regions (i.e. an additional $\approx 35\%$, when compared to the gaseous component; Mesa Delgado et al. 2009; Peimbert & Peimbert 2010; Espíritu et al. 2017).

Here we study the discrepancy between TIM and DM abundances from a different perspective. We will compute chemical evolution models adopting the initial mass function (IMF) by Kroupa (2002) with different M_{up} values, where M_{up} is the upper mass limit of the IMF. The M_{up} value is not the maximum stellar mass present in a given H II region, but the maximum mass of the IMF averaged over the age of the Galaxy.

In order to be able to study in depth as many details as possible, in this paper we will do a deeper study for a few of H II regions only. We selected the Orion Nebula because it is, by far, the most studied Galactic H II region. For our second object we selected M17 because it is the second most studied Galactic H II region, it is relatively close by, yet it has an appreciably different galactocentric distance; also it is a high-ionization H II region and thus we need not worry about the possible presence of neutral helium. Our last object is M8; it is one of the most studied Galactic H II regions, it has approximately the same galactocentric radius as M17.

Chemical evolution models for the MW have been built to reproduce robust observational constraints of the Galaxy. Some authors (e.g., Prantzos et al. 2018; Romano et al. 2010) constrain their CEMs by trying to fit the solar abundances at the Sun's age; other authors (e.g., Spitoni et al. 2019; Mollá et al. 2015) constrain their models by trying to reproduce the $[\alpha/\text{Fe}]$ - [Fe/H] trend shown by stars at the solar vicinity. Most of these models reproduce the current slope of the α/H gradients, but the predicted absolute values are different and, consequently, the predicted α enrichment efficiencies are different during the last few Gyrs of the evolution.

Unfortunately, the chemical gradients of H II regions cannot be used as solid constraints because, while the slope of the chemical gradient is widely accepted, the absolute values of H II region gradients found in the literature present a large dispersion; this becomes more pronounced when combined with gradients derived from other young objects.

To improve the quality of the models, it is important to fit the absolute value of the element abundances at the present time, not only the slope. These absolute values become critical to place restrictions



DM

4

12 16 20

8

R (kpc)

Fig. 1. Values of O/H as a function of the distance to the galactic center. Models for $t^2 = 0.00$ and for observed t^2 values versus two chemical evolution models with different M_{up} values. The filled and dashed lines represent the radial distribution obtained with the TIM $(M_{up} = 80M_{\odot})$ and DM $(M_{up} = 40M_{\odot})$ models, respectively. The circles represent the values derived from observations: TIM values based on recombination lines (black), TIM values based on the calibration by Peña-Guerrero et al. 2012 (empty), and DM values (grey). For further discussion, see Carigi et al. (2019a). The color figure can be viewed online.

16

20

on the CEMs. For example: since O, and other α elements, are mainly produced by massive stars, the absolute values of the gradient are critical to determine the M_{up} value of the IMF.

For simplicity we will use H, He, C, and O to represent the abundances by number, and X, Y, C, and O to represent the abundances by mass of these elements; Z represents the total heavy element abundance by mass.

2. H II REGIONS ABUNDANCES AND CHEMICAL EVOLUTION MODELS

2.1. O/H vs. Distance to the Galactic Center

In Figure 1 we present two sets of H II region data for O/H, one based on the TIM and the other based on the DM, as well as the best fit models to each data set. The observational data were compiled from Esteban et al. (2004, 2013, 2016, 2017); Fernández-Martín et al. (2017); García-Rojas et al. (2004); García-Rojas et al. (2005, 2006); García-Rojas et al. (2007); García-Rojas, Simón-Díaz & Esteban (2014). The models, built to reproduce the O/H gradient, come from Carigi et al. (2019a). For a more detailed description of the data selection, the model, and the fit see Carigi et al.(2019a). All abundances have been corrected by the fraction of O trapped in dust grains (Peimbert & Peimbert 2010; Peña-Guerrero et al. 2012; Espíritu et al. 2017). The figure illustrates very well that the curves required to fit the TIM and the DM data are quite different; it is thus no surprise that the M_{up} required by the models to fit the TIM data and the DM data are very different: while the M_{up} used to fit the TIM data amounts to 80 M_{\odot} , the value used to fit the DM data amounts to 40 M_{\odot} .

When comparing these values with the ones observed for young objects (B-stars, Cepheids), we find that the model based on the TIM values produces an excellent fit between 5 and 17 kpc, while the model based on the DM values fails to reproduce the observations (Carigi, Peimbert, & Peimbert 2019b).

2.2. Representation of the Chemical Evolution Models

We compute a set of nine chemical evolution models (CEMs) for MW like galaxies based on the work by Carigi et al. (2019a); these models differ only in the adopted M_{up} value. We present the output of these models for 6.2 kpc, corresponding to the average galactocentric distance of M17 (6.1 kpc) and M8 (6.3 kpc), and for 8.34 kpc, corresponding to the galactocentric distance of the Orion Nebula.

The initial abundances of our models are: X(0) = 0.7549, Y(0) = 0.2451, C(0) = 0, and O(0) = 0, where $Y(0) = 0.2451 \pm 0.0026$ is the primordial helium abundance derived by Valerdi et al. (2019). Their Y(0) result is in good agreement with the value derived by Planck Collaboration (2018) that amounts to $Y(0) = 0.24687 \pm 0.00076$.

We explore the M_{up} effects on the predicted value of Y, C, and O during the whole evolution. The galaxies are formed in an inside-out scenario of primordial infall, with the halo component formed from 0 to 1 Gyr and the disk component formed from 1 to 13 Gyr.

In Figures 2 and 3 we present the ΔO vs ΔY evolution for R = 6.2 and 8.34 kpc, respectively. Moreover, in Figures 4 and 5 we show ΔO vs ΔC evolution curves computed for the same radii. The evolution curves are presented for nine CEMs that consider $M_{up} = 25, 30, 35, 40, 50, 60, 80, 100, 150 M_{\odot}$. For comparison, we present observational data for M17, M8, and Orion, using the DM and the TIM.

In Carigi et al. (2019a), models were built to reproduce the radial behavior of the total O/H from 21 H II regions. To fit a representative absolute value of the gradient, they inferred two M_{up} values: one if the gaseous O/H values were determined

9.0

8.8

12 + log(O/H) .8 .8 .9

8.2

8.0

0 4 8 12

TIM

R (kpc)

Fig. 2. Chemical evolution for Y and O at a galactocentric distance of 6.2 kpc (approximatelly the distance of M17 and M8). The curves cover the entire evolution from the beginning (0 Gyr) to the present time (13 Gyr, magenta points), and each curve corresponds to a model with a different M_{up} . The squares represent the O and He abundances derived for M17 using the DM (grey) and the TIM (black); the diamonds represent the abundances derived for M8 using the DM (grey) and the TIM (black). The dotted magenta line connects the present-time values predicted by the models. Note that, the observed values should be compared with this magenta line to choose the better M_{up} values. The color figure can be viewed online.

from the DM $(M_{up} = 40 \ M_{\odot})$ and the other if the gaseous O/H values were determined from the TIM $(M_{up} = 80 \ M_{\odot})$, see Figure 1.

Also, for NGC 6822 (an irregular galaxy), Hernández-Martínez et al. (2011) built chemical evolution models to reproduce O/H values determined from DM and TIM, and obtained $M_{up} = 40 M_{\odot}$ and $M_{up} = 80 M_{\odot}$, respectively, in agreement with the values found for the MW.

In this work, we will obtain uncertainty bars for the M_{up} values, comparing the present-time abundances computed by models using different M_{up} values, with the O/H, He/H, and C/H abundances (and their error bars) for M17, M8, and Orion.

2.3. Object Selection

The approach in this study is to put quality over quantity; thus, we only use three objects: the Orion Nebula, M17, and M8, since they are the most studied Galactic H II regions. One very important characteristic of the Orion Nebula and M17 is that they are very bright (hence their many studies); as a consequence of this they have arguably the best O/H abundance ratio determinations. Another reason

Fig. 3. Chemical evolution for Y and O at a galactocentric distance of 8.34 kpc (Orion). The curves cover the entire evolution from the beginning (0 Gyr) to the present time (13 Gyr, magenta points), and each curve corresponds to a model with a different M_{up} . The circles represent the O and He abundances derived for Orion using the DM (grey) and the TIM (black). Note that the Orion data should be compared with the magenta line (theoretical present-time values). The color figure can be viewed online.

to select them is that they have noticeable different galactocentric radii. While the Orion Nebula is, by far, the best studied HII region, M17 has the advantage of being a high ionization H II region and thus we do not need to worry about an uncertain ICF(He). Unfortunately, there are no UV observations of M17 (probably due to its relatively high $c(H\beta) = 1.17$) and it is therefore not possible to derive C abundances using CELs and the direct method. We selected M8 because it is also nearby and very bright, its galactocentric distance is very similar to that of M17 (allowing us to present both of them using the same simulation and figures), it is the 4th most observed Galactic HII region, and probably the 3rd best suited for a study such as the one we present.

Tables 1, 2, and 3 show the total abundances (gas + dust) by mass derived for M17, M8, and Orion, respectively. The values were derived by transforming the abundances by number obtained by García-Rojas et al. (2007); Esteban et al. (2005) for O/H, He/H, and C/H and assuming that O is approximately 45% of Z. The first column of these tables shows the abundances derived through the DM (from CELs) and assumes a constant temperature over the observed volume, whereas the second column shows the abundances derived through the TIM (from RLs).





Time (Gyr

 $0.73 {\pm} 0.10$

M17: OBSERVED ΔY , O , AND C VALUES ^a				
	DM	TIM		
$\Delta Y \left(10^{-3} \right)$	$37.7 {\pm} 4.2$	30.1 ± 4.1		
$O(10^{-3})$	$4.81 {\pm} 0.48$	$8.56{\pm}0.86$		
$C(10^{-3})$		$6.27 {\pm} 0.63$		
$\Delta Y / \Delta O$	$7.86{\pm}1.39$	$3.52 {\pm} 0.72$		

TABLE 1

^aObservations from García-Rojas et al. (2007).

 $\Delta C / \Delta O$

TABLE 2

M8: OBSERVED ΔY , O, AND C VALUES^a

	DM	TIM
$\Delta Y \left(10^{-3} \right)$	42.2 ± 8.5	22.1 ± 7.8
$O(10^{-3})$	$4.80 {\pm} 0.49$	$7.42{\pm}0.79$
$C(10^{-3})$	$1.50\substack{+0.30\\-0.75}$	$5.30^{+1.07}_{-2.65}$
$\Delta Y / \Delta O$	$8.80{\pm}1.99$	$2.84{\pm}1.06$
$\Delta C/\Delta O$	$0.31\substack{+0.07 \\ -0.16}$	$0.68\substack{+0.15 \\ -0.35}$

^aObservations from García-Rojas et al. (2007).

As mentioned above, the line of sight in the direction of M17 has a relatively high reddening, and the $\lambda\lambda$ 1906-1909 Å [C III] lines are too obscured to have been observed. One might be interested in using the λ 4267 Å C II line to complete the DM determination; but λ 4267 Å C II cannot be used as part as the DM for the same reasons that the λ 4650 Å O II multiplet cannot be used (since it corresponds to the TIM). Therefore, while widely used, it should not be considered as part of the DM (and one should beware of authors that use it as part of the DM without a clear and consistent explanation on the ADF origin and its consequences).

2.4. O/H vs. He/H

Gaseous O/H abundances are readily available from many observational sources. However, they are frequently not converted to the total ISM O/H; to do this, it is necessary to include the fraction of O trapped in dust grains. It is estimated that this correction is between 0.07 and 0.13 dex for most H II regions (Mesa Delgado et al. 2009; Peimbert & Peimbert 2010; Espíritu et al. 2017); the exact value depends on the metallicity and on the efficiency of the dust destruction present within each H II region. Here we will include a correction of 0.12 dex for the Orion Nebula (Mesa Delgado et al.

TABLE 3

ORION	NEBULA:	OBSERVED	$\Delta Y, O, AND$
	C	VALUES ^a	

	DM	TIM
$\Delta Y \left(10^{-3} \right)$	33.2 ± 7.1	30.3 ± 7.0
$O(10^{-3})$	$4.92 {\pm} 0.45$	$7.39 {\pm} 0.60$
$C(10^{-3})$	$1.35\substack{+0.55\\-0.40}$	$2.81^{+0.42}_{-0.32}$
$\Delta Y / \Delta O$	$6.75 {\pm} 1.45$	$4.10{\pm}1.01$
$\Delta C / \Delta O$	$0.27^{+0.12}_{-0.08}$	$0.38^{+0.06}_{-0.05}$

^aObservations from Esteban et al. (2005).

2009; Espíritu et al. 2017) and 0.11 dex for both M17 and M8(Peimbert & Peimbert 2010).

The C/H abundance should also be corrected for dust depletion; this correction is expected to be similar or slightly smaller than the O/H correction (Esteban et al. 1998, 2009). Here we will assume a correction of 0.10 dex for all three H II regions.

Although most elements should include a correction due to dust depletion, the fact that He is an inert noble gas means that no correction will be necessary for the He/H abundances.

2.4.1. O vs. Y for M17 and M8

In Figure 2 we present the theoretical evolution of ΔO and ΔY vs time, for R = 6.2 kpc. We plot nine curves that correspond to the nine M_{up} values listed in § 2.2. The curves begin at t = 0 Gyr $(\Delta O = \Delta Y = 0)$ and end at 13 Gyr. We include the ΔO and ΔY values for M17 and M8, determined from the DM and the TIM (see Tables 1 and 2). In order to choose the M_{up} values that best reproduce the observational data, the top of each curve (the predicted values at present time (shown in magenta points) should be compared with the M17 and M8 data (the observed abundances for the ISM).

The curves evolve more rapidly to the right with increasing M_{up} values, because the O production for high mass stars (HMS) increases with the stellar mass. At 1 Gyr, when the halo formation ends, the curves present a loop due to the dilution of the ISM with primordial infall (Y = 0.2451, O = 0.0 Valerdi et al. 2019) that forms the disk. The rest of the evolution depends on the lifetime and the initial metallicity (Z) of the HMS and low-and-intermediate mass stars (LIMS), see Carigi & Peimbert (2008) and Carigi & Peimbert (2011).

Current Y values are almost constant for $M_{up} < 50 M_{\odot}$ and increase for $M_{up} > 50 M_{\odot}$ (see Table 4), because: (i) for low Z (equivalent to low O),

TABLE 4

	$\Delta Y (10^{-3})$	$O(10^{-3})$	$C(10^{-3})$	$\Delta Y / \Delta O$	$\Delta Y / \Delta C$
$\frac{111}{150}$	39.16	8.63	7.54	4.54	1.14
100	37.14	8.15	7.08	4.56	1.15
80	35.92	7.72	6.67	4.65	1.16
60	34.79	7.08	6.01	4.91	1.18
50	34.40	6.57	5.47	5.24	1.20
40	34.18	5.83	4.86	5.86	1.20
35	34.09	5.24	4.50	6.51	1.16
30	34.19	4.43	4.28	7.72	1.04
25	34.19	3.36	4.24	8.06	0.79

PRESENT DAY VALUES IN THE ISM PREDICTED BY THE MODELS FOR $R=6.2~{\rm KPC}$ (M17 AND M8)

HMS in the 8-25 M_{\odot} range are much more efficient to produce He than those in the 25-150 M_{\odot} range, (ii) for high Z, HMS in the 50-150 M_{\odot} range are very efficient to produce He, and (iii) for LIMS, the He contribution is not strongly Z-dependent.

A peculiarity of Figure 2 is the shape of the $M_{up} = 150 \ M_{\odot}$ curve, where the oxygen abundance diminishes between 7.8 and 10.0 Gyrs. This O dilution is caused by the huge amount of C ejected by very massive stars of high Z (see the description of the carbon evolution in § 2.5.1).

When comparing the observed O and He values for M17 with those derived from our models we find: (i) for the ΔO value determined with the DM, an IMF with a galactic M_{up} of 30 - 36 M_{\odot} , while (ii) for the ΔO value determined with the TIM, an $M_{up} > 75 M_{\odot}$, (iii) for the ΔY value determined with the DM, all values of M_{up} are allowed (the ΔY value is not very restrictive), and (iv) for the ΔY value determined with the TIM, an $M_{up} < 70 M_{\odot}(1\sigma)$ (allowing for all possible values at the 2σ level).

From the M8 values we find: (i) the ΔO value determined with the DM is nearly identical to the one determined for M17; therefore, the range determined for M_{up} is also of 30 - 36 M_{\odot} , (ii) the ΔO value determined with the TIM suggests M_{up} in the 52 - 120 M_{\odot} range, (iii) for the ΔY value determined with the DM, all values of M_{up} are allowed (the ΔY value is not very restrictive), and (iv) for the ΔY value determined with the TIM there is no solution at the 1 σ level, yet at the 2 σ level all solutions are allowed; this shows both the uncertainty of the ICF(He) and the lack of restriction produced by ΔY .

2.4.2. O vs. Y for the Orion Nebula

In Figure 3 we present curves of the theoretical evolution of ΔO and ΔY vs time for R = 8.34 kpc, corresponding to nine M_{up} values listed in § 2.2. Moreover, we include the ΔO and ΔY values for Orion, determined from the DM and the TIM (see Table 3). As in Figure 2, the top of each curve (in magenta) corresponds to the end of evolution (i.e. the present time; see Columns 2 and 3 of Table 5), and should be compared with the Orion data to choose the M_{up} values that best reproduce the observations.

In Figure 3, for any given M_{up} , the evolutionary curves reach lower ΔO and ΔY values than the corresponding coeval values in Figure 2, because the O/H gradient is negative for all MW-like models. Therefore, for any given time, the ΔO value (and Z value) for R = 8.34 kpc is lower than the ΔO value for R = 6.2 kpc. Consequently, very few massive stars of high Z form and, since they are more efficient He producers, the reached Y values are lower. In this figure, for the $M_{up} = 150$ curve, the O dilution is lower, due to the relative lack of massive stars of high Z (highly-efficient C producers).

When comparing the observed O and He values for Orion with those derived from our models, we find: (i) for the ΔO value determined with the DM, an M_{up} of 38 - 50 M_{\odot} ; (ii) for the ΔO value determined with the TIM, an M_{up} of 75 - 130 M_{\odot} , (iii) for the ΔY value determined with the DM, an $M_{up} > 35 M_{\odot}(1\sigma)$ (allowing for all possible values at the 1.2 σ level), and (iv) for the ΔY value from the TIM, again, all values of M_{up} are allowed.

2.5. O vs. C

2.5.1. O vs. C for M17 and M8

In Figure 4 we show the curves of theoretical evolution of ΔO and ΔC vs time, for R = 6.2 kpc, obTABLE 5

(ORION NEBULA) (ORION NEBULA)					
M_{up}	$\Delta Y(10^{-3})$	$O(10^{-3})$	$C(10^{-3})$	$\Delta Y / \Delta O$	$\Delta Y / \Delta C$
150	28.82	8.32	5.22	3.46	5.52
100	27.87	7.65	4.84	3.64	5.76
80	27.30	7.03	4.54	3.88	6.01
60	26.77	6.12	4.10	4.37	6.53
50	26.45	5.45	3.77	4.85	7.02
40	26.12	4.60	3.42	5.68	7.64
35	25.91	4.02	3.24	6.45	8.00
30	25.50	3.30	3.11	7.73	8.20
25	25.24	2.45	3.04	10.30	8.30

DESENT DAV VALUES IN THE ISM DEFICTED DV THE MODELS FOR D = 9.24 KDC



Fig. 4. Chemical evolution for C and O at a galactocentric distance of 6.2 kpc (approximately the distance of M17 and M8). The curves and points are similar to those in Figure 2. Notice, however, that there are only 3 points in this figure instead of the 4 in Figure 2: unfortunately there is no restriction on C for the DM of M17 and, instead of the fourth point the shaded vertical band is the ΔO predicted by the DM. Again, the observed values should be compared with the magenta curve. The color figure can be viewed online.

tained from our nine CEMs. The current values, at the top end of the curves, are presented in Columns 3 and 4 of Table 4. We include the observed ΔO and ΔC values for M17 and M8 determined from the TIM. The gray diamond represents the ΔO and ΔC values determined for M8, while the shaded vertical bar represents the ΔO for M17 using the DM combined with the lack of a C determination available from the DM (see Table 1).

A remark on the ΔC determination for M8: according to the recent ICFs computed for giant H II regions (Amayo, Delgado-Inglada, Stasińska, 2020, in prep.), the use of $C/O = C^{++}/O^{++}$ in M8 may underestimate the real C/O value by up to ≈ 0.3 dex. However, these ICFs may not be adequate for Galactic H II regions where only a small area is observed. We decided not to change the value of C/H but to increase the associated error bars.

Current C values are almost constant for $M_{up} \lesssim 35 M_{\odot}$ and increase significantly for $M_{up} \gtrsim 40 M_{\odot}$ (see Table 4); this happens for two main reasons: (i) stars in the 40-150 M_{\odot} range produce much more C when they are more metal rich, and (ii) stars in the 1-3 M_{\odot} range produce more C when they are metal poor. Due to the the LMS enrichment contribution, HMS of high Zcontribute at times similar than LMS of low Z (e.g. Akerman et al. 2004; Carigi et al. 2005; Carigi & Peimbert 2011).

When comparing the observed ΔC value for M17 determined with the TIM with those derived from the CEMs, we find an M_{up} in the 55 - 95 M_{\odot} range. Using the determination from the M8's ΔC measurements obtained with the TIM, we only find an upper limit $M_{up} < 72 \ M_{\odot}$; the lack of a lower limit is due to the large uncertainty on the lower limit of the ICF. When determining a ΔC value for M8 using the DM, the value is not compatible with our models, as the DM falls short of our models by a factor of at least 2.5 (for the lowest $M_{up} = 30 M_{\odot}$), and probably by a factor of about 3.5 or more (for a more reasonable $M_{up} \gtrsim 70 M_{\odot}$). Finally, regarding the DM determination for M17: due to the high reddening of M17, it has not been observed in the UV, and therefore it has not been possible to obtain [C III] intensities, nor to determine DM abundances. The $\Delta O M_{up}$ determinations are the same as those in Figure 2.

12.0

10.5

9.0

7.5

6.0

4.5

3.0

1.5

0.0

10

Time (Gyr

Fig. 5. Chemical evolution for C and O for a galactocentric distance of 8.34 kpc (Orion). The curves and points are similar to those in Figure 3. The Orion data should be compared with the magenta line. The color figure can be viewed online.

=25, 30, 35, 40, 50, 60, 80, 100, 150 M

-3)

4

 $\Delta O (10^{-1})$

8

2.5.2. O vs. C for Orion

In Figure 5 we show the theoretical evolution of ΔO and ΔC vs time for R = 8.34 kpc, for each of our nine M_{up} values. Moreover, we include the ΔO and ΔC values for Orion, determined from the DM and the TIM (see Table 3). In this figure, for any given M_{up} , the evolutionary curves reach lower O and Cvalues than the coeval values in Figure 4, because the O/H and C/H gradients are negative for all the MW-like models.

By comparing the observed C values for Orion with the current values derived from our models (see Table 5), we find that neither set of observed ΔC abundances (neither DM nor TIM) are consistent with the theoretical predictions. Since the ΔO values and determinations are the same as those in Figure 3 it seems that the C abundance in Orion is lower than expected; however, the C/H abundance in the Orion nebula is noticeable smaller (about 0.1 dex) than in NGC 3603 ($R = 8.65 \,\mathrm{kpc}$) and NGC 3576 ($R = 7.46 \,\mathrm{kpc}$), which are the two H II regions with a galactocentric distance closest to the one of the Orion nebula (Esteban et al. 2004; García-Rojas et al. 2004; García-Rojas et al. 2006). There are two causes that may explain this low value. The first one is the use of an inadequate ionization correction factor (ICF); but according to a recent study on ICFs for giant H II regions (Amayo et al. (2020) in prep.) it is adequate to use $C/O = C^{++}/O^{++}$ to compute C abundances. The second one is the possibility of this nebula having more C atoms deposited in dust grains and thus, a lower gaseous abundance of C; a clue that the dust properties in Orion are different than in other HII regions is the high total to selective absorption ratio present in Orion, that amounts to about $R_V = E_V/E_{B-V} \approx 5.5$ (Peimbert & Costero 1969; Esteban et al. 2004), while for most other objects it amounts to about $R_V = 3.1$ (Cardelli et al. 1989). If we consider that Orion could have a slightly higher (0.1 dex) total C abundance, we find, for the TIM, the total gas plus dust ratio of C/H for the Orion nebula to be $12 + \log(C/H) \approx 8.53 \pm 0.08$; this value represents a slightly lower M_{up} than the one derived from O/H. On the other hand, we find, for the DM, the total C/H to be $12 + \log(C/H) \approx 8.22 \pm 0.10$, still 0.1 dex less than our lowest model, approximately 0.2 dex lower than the $M_{up} \approx 40 \ M_{\odot}$ favored by the O/H determinations of the DM, and approximately 0.3 dex lower than what is required to be consistent with the more favored $M_{up} \lesssim 80 \ M_{\odot}$ values.

3. THE CHEMISTRY HAS A BETTER MEMORY THAN THE LIGHT

The lifetime of massive stars (few Myr) is very short compared to the age of galaxies (several Gyr); consequently, massive stars are not an important component of the light of the majority of the observed galaxies; yet, when massive stars die, their contributions are quickly incorporated into the chemistry of the ISM.

The chemical composition of an H II region is the result of the whole history of the chemical evolution of any given galaxy; therefore the chemical composition of H II regions can be compared with estimates of the present day chemistry derived from galactic CEMs; and we can infer the amount of formed massive stars (equivalently, the M_{up} value), comparing the chemical abundances in the ISM with those obtained from CEMs.

To study the most massive stars in the MW by looking for them at present has several major inconveniences: there are very few of these stars, this is compounded by the fact that they must be formed in very massive giant molecular clouds (GMCs), and they are the first ones to evolve (first shedding mass in strong stellar winds, and then going supernova, all this before the GMC is dissipated by the combined effect of the stars that are evolving inside it). Thus, they are usually obscured and very difficult to observe during their short lifespan. Moreover, most massive stars that can be observed today may not be representative of the most massive stars that have existed during the evolution of the Galaxy, i.e. the stars that have contributed to the evolution of the



6

1

0

Orion TIM Orion DM

Mup

0

2

R=8.34 kpc

chemistry of the present day ISM as well as the chemistry available during the most recent star formation.

Based on the ΔY comparison, between the observed values (from TIM or DM) and the preset-time predicted values (from CEMs), we cannot exclude any M_{up} value in the 25 - 150 M_{\odot} range. However, from the ΔO comparison, we find agreement for M_{up} values in the $30 < M_{up} < 50 M_{\odot}$ range for the DM, and in the $75 < M_{up} < 120 M_{\odot}$ range for the TIM. Moreover, from the ΔC comparison, we find agreement in the $55 < M_{up} < 95 M_{\odot}$ range for the TIM, and a suggestion of a very small M_{up} for the DM $(M_{up} < 25 M_{\odot})$. Therefore, by comparing TIM with CEMs, the best M_{up} values are in the $70 \leq M_{up} \leq 100 M_{\odot}$ range; and by comparing DM with CEMs, the best M_{up} values are in the $25 \leq M_{up} \leq 45 M_{\odot}$ range.

Weidner et al. (2013) in their Figure 3 showed the dependence of the star formation rate (SFR) on the integrated galactic stellar initial mass function (IGIMF, called IMF in our CEMs) for different power-law indexes, α , (for initial stellar masses between 1.3 M_{\odot} and M_{up}). They noted that IMF for a SFR $\approx 1 \ M_{\odot}$ /yr corresponds to an $\alpha = 2.6$ and to an $M_{up} \approx 100 \ M_{\odot}$; and that these values are in agreement with the MW (we use $\alpha = 2.7$ for our CEMs). On the other hand, the SFR required to obtain $M_{up} \approx 40 \ M_{\odot}$ is SFR $\approx 10^{-2} \ M_{\odot}$ /yr (as well as an $\alpha \approx 2.8$).

Regarding the observational determinations of the SFR: (i) the MW, a spiral galaxy (Sbc) with total stellar mass $\approx 10^{11} M_{\odot}$, presents a galaxywide SFR $\approx 0.7 - 2.3 M_{\odot}/\text{yr}$ (Robitaille & Whitney 2010; Chomiuk & Povich 2011); (ii) NGC 300, a small spiral galaxy (Sd) with total stellar mass $= 1.9 \times 10^9 M_{\odot}$, presents a galaxy-wide SFR $= 0.08 - 0.30 M_{\odot}/\text{yr}$, approximately an order of magnitude smaller than the SFR in the MW (e.g. Kang et al. 2016, and references therein).

Moreover, based on spatially-resolved spectroscopic properties of low-redshift star-forming galaxies, Sánchez (2019) showed in his Figure 7 a difference of approximately 1 order of magnitude between the SFR of Sbc galaxies with stellar mass $\approx 10^{11}$ M_{\odot} (as the MW galaxy) and the SFR of Sd galaxies with stellar mass $\approx 2 \times 10^9 M_{\odot}$ (as NGC 300).

Therefore, the M_{up} values we derive from the TIM are consistent with the SFR of the MW galaxy, while the M_{up} values derived from the DM are consistent with a galaxy with a mass and SFR similar to those of NGC 300, but not with the mass and SFR of the MW.

Moreover the abundances derived from the TIM are consistent with those derived from observation of other young objects in the MW (Cepheids, B stars), while the abundances derived from the DM are approximately 0.25 dex too small.

The chemical composition of a given H II region is the result of the evolution of the ISM throughout the history of our Galaxy. Therefore, the chemical abundances of an H II region do not depend on the IMF of the observed H II region. In particular the most massive star of a given H II region is not representative of the most massive stars of the galactic IMF.

4. CONCLUSIONS

We have computed nine chemical evolution models (CEM) of a MW like galaxy; the only difference among these models is the IMF, specifically its $M_{\mu\nu}$ value, that ranges between 25 and 150 M_{\odot} . We compare the model predictions with the O/H, He/H, and C/H values derived for three Galactic H II regions: M17, M8, and Orion. We compute the abundances by two different methods: the DM (direct method) and the TIM (temperature independent method); these methods have always given different results. We selected these objects because they probably have the best O/H determinations, and because their galactocentric radii are different enough to be useful as independent constraints for the CEMs (6.2 and)8.34 kpc, respectively). The comparison between models and observations tells us which is the $M_{\mu\nu}$ that better fits each set of observations.

It is useful to remember that the chemistry has a better memory than the observed UV light. In other words: the chemistry will explore the average M_{up} over the lifespan of the MW, while any measurement of the UV radiation or of the most massive stars observed can only be a reflection of the present day M_{up} (and can potentially have significant biases toward lower masses).

When comparing the models with the DM abundances we find: for ΔO , a $30 < M_{up} < 50 M_{\odot}$ range; for ΔC , $M_{up} < 25 M_{\odot}$, while Y does not provide a significant restriction. Overall, the DM produces preferred values in the $25 < M_{up} < 45 M_{\odot}$ range. On the other hand, when using the TIM abundances we find: for ΔO , a $52 < M_{up} < 150 M_{\odot}$ range; for ΔC , a $25 < M_{up} < 95 M_{\odot}$ range, while Y suggests a smaller value, but does not provide a significant restriction. Overall the TIM produces preferred values in the $70 < M_{up} < 100 M_{\odot}$ range.

Moreover the M_{up} in a given galaxy is directly related to the SFR, and the SFR is directly related to the mass of any given galaxy. A MW like galaxy, with a SFR $\approx 1 \ M_{\odot}/\text{yr}$, is expected to have an $M_{up} \approx 100 \ M_{\odot}$, in good agreement with the TIM determination, but not with the DM determination (which would be more consistent with an Sd galaxy with a SFR $\approx 0.01 \ M_{\odot}/\text{yr}$).

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REFERENCES

- Akerman, C. J., Carigi, L., Nissen, P. E., et al. 2004, A&A, 414, 931
- Carigi, L., & Peimbert, M. 2008, RMxAA, 44, 341
- _____. 2011, RMxAA, 47, 139
- Carigi, L., Peimbert, M., Esteban, C., & García-Rojas, J. 2005, ApJ, 623, 213
- Carigi, L., Peimbert, M., & Peimbert, A. 2019a, ApJ, 873, 107
 - _____. 2019b, arXiv:1906.07804
- Cardelli, J. A., Clayton, G. C., & Mathis, J. S. 1989, ApJ, 345, 245
- Chomiuk, L., & Povich, M. S. 2011, AJ, 142, 197
- Espíritu, J. N., Peimbert, A., Delgado-Inglada, G., & Ruiz M. T. 2017, RMxAA, 53, 95
- Esteban, C., Bresolin, F., Peimbert, M., et al. 2009, ApJ, 700, 654
- Esteban, C., Carigi, L., Copetti, M. V. F., et al. 2013, MNRAS, 433, 382
- Esteban, C., Fang, X., García-Rojas, J., & Toribio San Cipriano, L. 2017, MNRAS, 471, 987
- Esteban, C., García-Rojas, J., Peimbert, M., et al. 2005, ApJ, 618, 95
- Esteban, C., Mesa-Delgado, A., Morisset, C., & García-Rojas, J. 2016, MNRAS, 460, 4038
- Esteban, C., Peimbert, M., García-Rojas, J., et al. 2004, MNRAS, 355, 229
- Esteban, C., Peimbert, M., Torres-Peimbert, S., & Escalante, V. 1998, MNRAS, 295, 401
- Esteban, C., Toribio San Cipriano, L., & García-Rojas, J. 2018, cagn.conf, 23
- Fernández-Martín, A., Pérez-Montero, E., Vílchez, J. M., & Mampaso, A. 2017, A&A, 597, 84
- García-Rojas, J. 2020, Reviews in "Frontiers of Modern Astrophysics: From Space Debris to Cosmology" (eds Kabath, Jones and Skarka; publisher Springer Nature) funded by the European Union Erasmus + Strategic Partnership grant "Per Aspera Ad Astra

Simul" 2017-1-CZ01-KA203-035562

- García-Rojas, J., Esteban, C., Peimbert, A., et al. 2005, MNRAS, 362, 301
- García-Rojas, J., Esteban, C., Peimbert, A. et al. 2007, RMxAA, 43, 3
- García-Rojas, J., Esteban, C., Peimbert, M. et al. 2004, ApJS, 153, 501
- García-Rojas, J., Esteban, C., Peimbert, M., et al. 2006, MNRAS, 368, 253
- García-Rojas, J., Simón-Díaz, S., & Esteban, C. 2014, A&A, 571, 93
- García-Rojas, J., Wesson, R., Boffin, H. M. J., et al. 2019, arXiv:190406763
- Hernández-Martínez, L., Carigi, L., Peña, M., & Peimbert, M. 2011, A&A, 535, 118
- Kennicutt, R. C. & Evans, N. J. 2012, ARA&A, 50, 531
- Kang, X., Zhang, F., Chang, R., Wang, L., & Cheng, L. 2016, A&A, 585, 20
- Kroupa, P. 2002, Sci, 295, 82
- Mesa Delgado, A., Esteban, C., García-Rojas, J., et al. 2009, MNRAS, 395, 855
- Mollá, M., Cavichia, O., Gavilán, M., & Gibson, B. K. 2015, MNRAS, 451, 3693
- Nicholls, D. C., Dopita, M. A., & Sutherland R. S. 2012, ApJ, 752, 148
- Peimbert, A. & Peimbert, M. 2010, ApJ, 724, 791
- Peimbert, M. 1967, ApJ, 150, 825

_____. 2019, arXiv: 1905.01244

- Peimbert, M. & Costero, R. 1969, BOTT, 5, 3
- Peimbert, M., Peimbert, A., & Delgado-Inglada, G. 2017, PASP, 129, 2001
- Peña-Guerrero, M. A., Peimbert, A., & Peimbert, M. 2012, ApJ, 756, 14
- Pérez-Montero, E. 2017, PASP, 129, 3001
- Planck Collaboration 2018, arXiv:1807.06209
- Prantzos, N., Abia, C., Limongi, M., Chieffi, A., & Cristallo, S. 2018 MNRAS, 476, 3432
- Robitaille, T. P. & Whitney, B. A. 2010, ApJ, 710, 11
- Rodríguez, M. & García-Rojas, J. 2010, ApJ, 708, 1551
- Romano, D., Karakas, A. I., Tosi, M., & Matteucci, F. 2010, A&A, 522, 32
- Sánchez, S. F. 2019, arXiv:1911.06925
- Spitoni, E., Silva Aguirre, V., Matteucci, F., Calura, F., & Grisoni, V. 2019, A&A, 623, 60
- Torres-Peimbert, S., Peimbert, M., & Peña, M. 1990, A&A, 233, 540
- Tsamis, Y. G., Barlow, M. J., Liu, X. W., Danziger, I. J., & Storey, P. J. 2003, MNRAS, 338, 687
- Valerdi, M., Peimbert, A., Peimbert, M., & Sixtos, A. 2019, ApJ, 876, 98
- Weidner, C., Kroupa, P., Pflamm-Altenburg, J., & Vazdekis, A. 2013, MNRAS, 436, 3309
- Carigi, L., Delgado-Inglada, G., Peimbert, A., and Peimbert, M.: Instituto de Astronomía, Universidad Nacional Autónoma de México, Apdo. Postal 70-264, México, C.P. 04510, CdMx, México (carigi, gdelgado, antonio, peimbert@astro.unam.mx).